

## Development of laser wakefield electron acceleration driven by high-repetition rate ultrafast lasers

M.-W. Lin

<sup>1</sup> Institute of Nuclear Engineering and Science, National Tsing Hua University  
e-mail (speaker): mwlin@mx.nthu.edu.tw

Raising the repetition rate of driving laser pulses can straightforwardly increase the average current of accelerated electrons from laser wakefield acceleration (LWFA), such that a high radiation yield can be achieved for developing advanced X-ray radiography and radiation therapy. To date, 200-mJ, 1030-nm pulses with a duration of 1.1 ps can be produced from a thin-disc Yb:YAG system at a repetition rate of 5 kHz [1]. By introducing these 1030-nm, ps-level pulses with a technique of post spectral broadening and pulse compression, such as the gas-filled multipass cell [2] or the multiple-plates continuum [3], pulses with a peak power  $< 4$  TW can be acquired to drive LWFA at kHz frequencies with the goal to increase the average electron current by 2-3 orders of magnitude higher than conventional LWFA  $\leq 10$  Hz.

In principle, a normalized vector potential  $a_0 \approx 2$  or a laser intensity  $I_L \gg 10^{18}$  W/cm<sup>2</sup> is required to drive LWFA in the 3D nonlinear regime [4]. In order for the few-TW LWFA to be successful, a gas target capable of providing a plasma density  $n_{ep} \gg 10^{19}$  cm<sup>-3</sup> and inducing a strong self-focusing of the pump pulse has to be invented. This self-focusing effect is characterized by the critical power  $P_{cr} \approx 17(n_{cr}/n_{ep})$  GW, depending on the plasma critical density  $n_{cr} = \epsilon_0 m_e \omega_L^2 / q_e^2$  and the laser frequency  $\omega_L$ . Our experiences indicate that  $P_L \sim 3 P_{cr}$  is favorable for conducting the few-TW LWFA. As the transverse radius  $w$  of the self-focused pump pulse gradually reduces to  $2w \sim \lambda_p$  (typically 3 - 5  $\mu$ m in operation), a greatly increased laser intensity is realized for driving plasma waves.

One of our approaches for realizing a stable source of tens-of-MeV electrons from LWFA is based on the recent experimental result achieved when  $\sim 3$  TW laser pulses are focused onto a pulsed thin gas jet of dense nitrogen. As shown in Fig. 1(a), the thin, high-density nitrogen gas jet was produced from the gas flow out of a 152- $\mu$ m diameter orifice under a high backing pressure. The nozzle consists of an orifice module mounted on a gas tube with the gas flow controlled by a pulsed valve underneath. Density of the gas atoms/plasma in the target region is varied by adjusting the backing pressure. The produced plasmas are measured by the probe of shadowgraphic images with a wavefront sensor, from which the plasma electron density distributions of the nitrogen target can be retrieved as shown in Fig. 1(b). When a 3.2-TW pulse is introduced with a backing pressure of 600 psi, the produced plasma density is fitted into a Gaussian profile with a width  $\sim 860$   $\mu$ m (in full width at half maximum, FWHM) and a peak density  $\sim 2.8 \times 10^{19}$  cm<sup>-3</sup>. The generated electron beam exhibited relatively small divergence properties for  $\theta_x \sim 20$  mrad

and  $\theta_y \sim 10$  mrad in FWHM as shown in Fig. 1(c). Under this condition, electrons were obtained with quasi-monoenergetic spectrum peaked at  $\sim 11$  MeV with an energy spread  $\sim 17$  MeV in FWHM and a bunch charge  $\sim 22$  pC ( $>5$  MeV) as shown in Fig. 1(d); more importantly, they are produced with satisfactory reproducibility under  $\sim 17\%$  energy and  $\sim 15\%$  charge stability. The experiments verified that such a scheme is a viable choice for generating stable tens-of-MeV electrons with few-TW laser pulses at a high-repetition rate.

### References

1. T. Nubbemeyer, M. Kaumanns, M. Ueffing, M. Gorjan, A. Alismail, H. Fattahi, J. Brons, O. Pronin, H. G. Barros, Z. Major, T. Metzger, D. Sutter, and F. Krausz, *Opt. Lett.* 42 1381 (2017).
2. M. Kaumanns, V. Pervak, D. Korman, V. Leshchenko, A. Kessel, M. Ueffing, Y. Chen, and T. Nubbemeyer, *Opt. Lett.* 43 5877 (2018).
3. C.-H. Lu, Y.-J. Tsou, H.-Y. Chen, B.-H. Chen, Y.-C. Cheng, S.-D. Yang, M.-C. Chen, C.-C. Hsu, and A. H. Kung, *Optica* 1, 400, (2014).
4. W. Lu, M. Tzoufras, C. Joshi, F. S. Tsung, W. B. Mori, J. Vieira, R. A. Fonseca, and L. O. Silva, *Phys. Rev. ST Accel. Beams* 13, 091301 (2010)

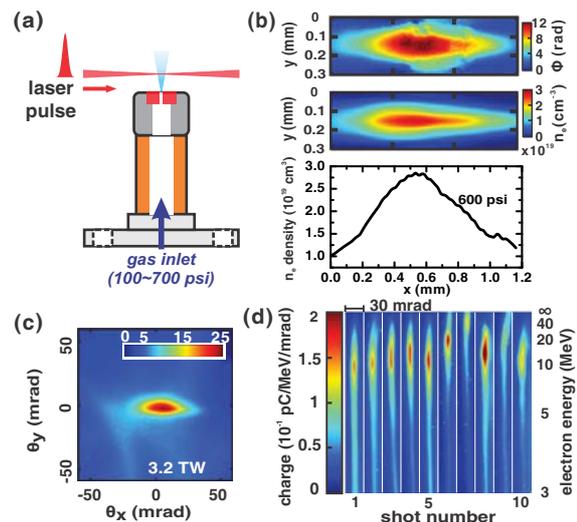


Figure 1 (a) Structure of the gas nozzle and the pump beam path. (b) The spatial phase shift of a shadowgraphic image, the retrieved plasma electron density distribution, and the density line profile of a nitrogen plasma produced by 3.2-TW laser power and 600-psi backing pressure. (c) Spatial profiles of accelerated electrons and (d) 10 consecutive images for the dispersed electrons generated with 3.2-TW pulses.