

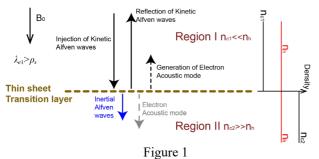
6th Asia-Pacific Conference on Plasma Physics, 9-14 Oct, 2022, Remote e-conference **Mode conversion from kinetic Alfvén waves to modified electron acoustic waves** Run Shi¹, Jun Liang²

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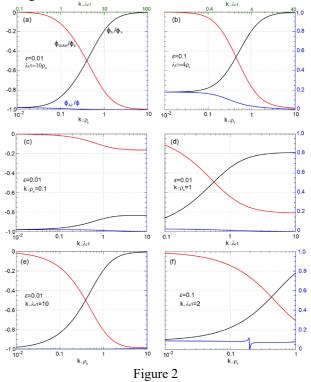
Possible mode conversion from kinetic Alfvén wave to modified electron acoustic wave is examined based on a multi-fluid model involving two electron populations. The mode conversion transpires when a kinetic Alfvén wave propagates through a transition between a hotelectron-dominant region and a cold-electron-dominant region.

We take advantage the fluid model in low-beta regime, and keep the parallel momentum equations of electrons with different temperatures



Region I denotes the region with dominantly hot electrons $(n_h \gg n_{c1})$, and Region II is the region with dominantly cold electrons $(n_h \ll n_{c2})$. The transition layer is approximated as a thin sheet. The hot electron density is assumed constant, while the density of cold electrons features a stepwise distribution across the transition boundary. The ambient magnetic field is assumed constant in both Region I and Region II.

The coefficients of the wave reflection ϕ_{1r}/ϕ_{1i} (black), transmission ϕ_{2IAW}/ϕ_{1i} (blue), and mode conversion ϕ_{1EAW}/ϕ_{1i} (red) are illustrated in Figure 2. Figures 2a and 2b show the dependence on k_{\perp} with fixed λ_{e1} , ρ_s , and $\varepsilon = \lambda_{e2}/\lambda_{e1}$. At smaller k_{\perp} , the electron acoustic mode plays a minor role. Both panels show that the mode conversion becomes significant with increasing k_{\perp} , while the KAW reflection becomes weaker. Figures2c exhibits the role of λ_{e1} with fixed k_{\perp} , ε and ρ_s . As expected, the magnitude of the mode conversion coefficient begins to increase as $k_{\perp}\lambda_{e1}$ approaches 1, and becomes stable when $k_{\perp}\lambda_{e1}$ >>1. Figure 2d is the same as Figure 2a except that ρ_s is 10 times bigger and hereby leads to greater variation with λ_{e1} . Similar to the effect of λ_{e1} , the increase of ρ_s (Figures 2e and 2f) also leads to stronger mode conversion.



References

R. L. Lysak and W. Lotko, J. Geophys. Res. 101, 5085 (1996). https://doi.org/10.1029/95JA03712
K. Stasiewicz, P. Bellan, C. C. Chaston, C. A. Kletzing, R. L. Lysak, J. Maggs, O. Pokhotelov, C. Seyler, P. Shukla, L. Stenflo et al., Space Sci. Rev. 92, 423 (2000). https://doi.org/10.1023/A:1005207202143