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## Study of ion kinetic effects on plasma distributions

by anisotropic-ion-pressure plasma fluid simulation

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Currently, efforts are being made to improve edge plasma transport simulations to establish methods for controlling the heat load on the divertor.

In plasma fluid models primarily used for scrape-off layer (SOL) plasma, the conductive heat flux along magnetic field lines is represented by the Spitzer-Härm (SH) heat flux  $q^{||SH|}$ [1]. However, the actual heat flux becomes smaller than the SH heat flux when the mean-free path is longer than the SOL characteristic length. This is referred to as the kinetic effect. In the SONIC code [2,3], a free streaming energy (FSE)limited model is employed for the kinetic effect. According to the SONIC simulation of JA DEMO [4], the kinetic effect of ion heat flux leads to remarkable changes in the SOL plasma profiles [3]. In these simulations, however, the anisotropy between the ion temperature parallel and perpendicular to the magnetic field direction, denoted as  $T_{i||}, T_{i\perp}$ , which might be significant in the low-collisionality regime [5], have not been considered.

In this study, we aim to study the ion kinetic effects in the presence of anisotropic ion pressure. A plasma fluid model based on the anisotropic ion pressure (AIP model) [6,7] is applied to a tandem mirror device GAMMA 10/PDX which exhibits significant ion pressure anisotropy [8]. In the AIP model, the ion conductive heat flux is evaluated based on a simple FSE-limited model [6].

$$q_{\sigma}^{||\text{FSE}} = \left[1 + \frac{\left|q_{\sigma}^{||\text{SH}}\right|}{\alpha_{\sigma}q_{\sigma}^{\text{FS}}}\right]^{-1} q_{\sigma}^{||\text{SH}} \left(\sigma \in \{i||, i \perp\}\right)$$
(1)

$$q_{\sigma}^{||\text{SH}} = -\kappa_{\sigma}^{||\text{SH}} (\partial T_{\sigma} / \partial s), \qquad q_{\sigma}^{||\text{FS}} = p_{\sigma} \sqrt{T_{\text{i}||} / m_{\text{i}}} \qquad (2)$$

Here,  $\kappa_{\sigma}^{||SH|}$  represents the SH thermal conductivity, *T* is the temperature, *p* is the pressure, *m*<sub>i</sub> is the ion mass, and *s* is the coordinate in the magnetic field direction. Equation (1) describes the FSE-limited model using the harmonic mean of  $q_{\sigma}^{||SH|}$  and  $q_{\sigma}^{||FS}$ , where  $q_{\sigma}^{||FS}$  is the free-streaming heat flux as a collisionless limit. The kinetic effects occur depending on collisionality, and the degree of these effects is defined by the heat flux limiter coefficients ( $\alpha_{\sigma}$ ). Note that, in the AIP model, the parallel ion conductive heat flux is divided into two components  $q_{i||}^{||}$  and  $q_{i\perp}^{||}$  which transport the parallel and perpendicular energy components in the parallel direction, respectively. In this study, we vary the value of  $\alpha_{\sigma}$  and change the kinetic effects in each component.

Figure 1 shows the spatial distribution of ion temperature anisotropy  $(T_{i\perp}/T_{i\parallel})$  in the GAMMA



**Figure 1.** Spatial distributions of magnetic field intensity *B* in GAMMA 10/PDX and ion temperature anisotropy  $T_{i\perp}/T_{i\parallel}$  for various  $\alpha_{i\parallel}$  values. The center (*s* = 0) is a symmetric point.

10/PDX device for various  $\alpha_{i|i}$  values with fixed  $\alpha_{i\perp} =$ 0.5. By the increase of  $\alpha_{i||}$ , both  $T_{i||}$  and  $T_{i\perp}$  in the central region (|s| < 2.8 m) are decreased. On the other hand, it is found that the anisotropy is unchanged  $(T_{i\perp}/T_{i\parallel} \approx 2.5)$ . Simulations are conducted in a uniform magnetic field system with the same collisionality. By increasing  $\alpha_{i||}$ , an increase in  $T_{i\perp}/T_{i||}$  is observed clearly. As a common phenomenon in both systems, increasing  $\alpha_{i|i}$  leads to an increase in  $q_{i|i}^{|i|}$  (> 0) and a resultant decrease in  $T_{i||}$ . Moreover, the nonlinear changes in the plasma distribution also affect  $q_{i1}^{||}$ . In the case of GAMMA 10/PDX, particle flow is hindered due to mirror confinement, and the conductive heat flux is comparable in magnitude to the convective heat flux. As a result, the decrease in  $q_{i\perp}^{||}$  (<0) brings about a decrease in  $T_{i\perp}$ . That is why the change in  $T_{i\perp}/T_{i\parallel}$  is minimal in GAMMA 10/PDX. In the presentation, the effect of  $\alpha_{i1}$  on the plasma distribution is also discussed.

References

[1] S.I. Braginskii, Reviews of Plasma Physics, vol.1

(Consultants Bureau, New York, 1965), 205.

[2] K. Shimizu et al., Nucl. Fusion 49 (2009) 065028.

[3] Y. Homma, Plasma Phys. Control. Fusion **64** (2022) 045020.

[4] Y. Sakamoto et al., DEMO concept development and assessment of relevant technologies [FIP/3-4Rb], 25<sup>th</sup> IAEA Int. Conf. on Fusion Energy (Saint Petersburg, 2014)

[5] A. Froese et al., Plasma Fusion Res. 5 (2010) 026

[6] S. Togo et al., Nucl. Fusion **59** (2019) 076041.

[7] S. Togo et al., J. Adv. Simulat. Sci. Eng.9 (2022) 185.

[8] M. Ichimura et al., Nucl Fusion 39 (1999) 1995.